

## Chapter 9 Center of Mass and Linear Momentum

- we can simplify the complicated motion of a system of objects by determining a special point of the system — the *center of mass* of that system.
- Although the motion of an object could be complicated, it has one special point – the center of mass – that does move in a simple parabolic path.

### The Center of Mass

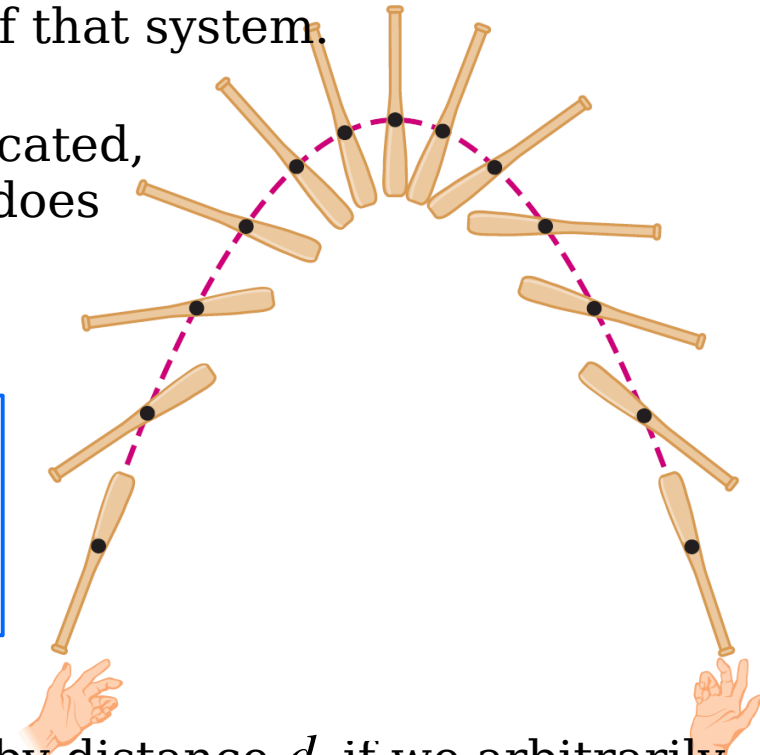
The **center of mass** of a system of particles is the point that moves as though (1) all of the system's mass were concentrated there and (2) all external forces were applied there.

### Systems of Particles

- For two particles of masses  $m_1$  and  $m_2$  separated by distance  $d$ , if we arbitrarily choose the origin of an  $x$  axis to coincide with the particle of  $m_1$ , then we *define* the position of the center of mass (com) of this two-particle system to be

$$x_{com} = \frac{m_2}{m_1 + m_2} d$$

- For example, if  $m_2 = 0$ , then  $x_{com} = 0$ ;  
if  $m_1 = 0$ , then  $x_{com} = d$ ;  
if  $m_1 = m_2$ , then  $x_{com} = d/2$ .



- the center of mass must lie somewhere between the 2 particles.

- For a more generalized situation, in which the coordinate system has been shifted, then the position of the center of mass is

$$x_{\text{com}} = \frac{m_1 x_1 + m_2 x_2}{m_1 + m_2} = \frac{m_1 x_1 + m_2 x_2}{M}$$

where  $M = m_1 + m_2$ , the total mass of that system.

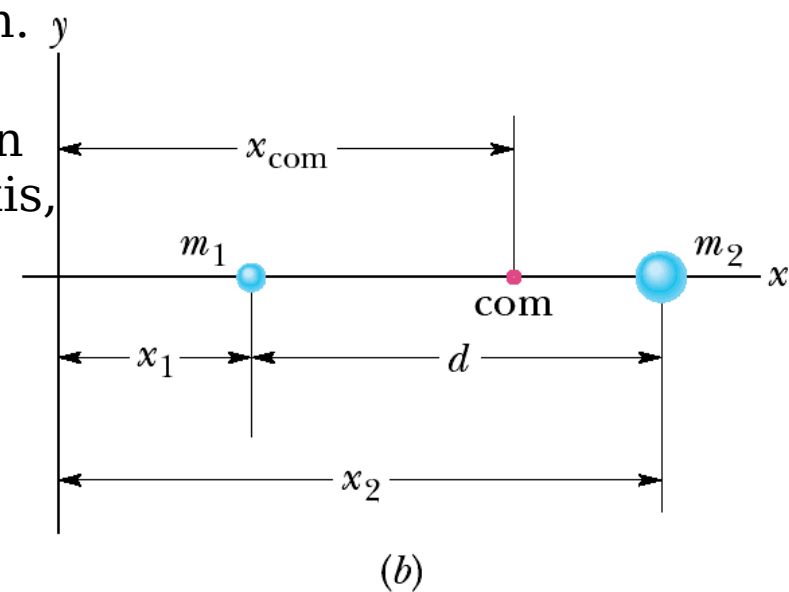
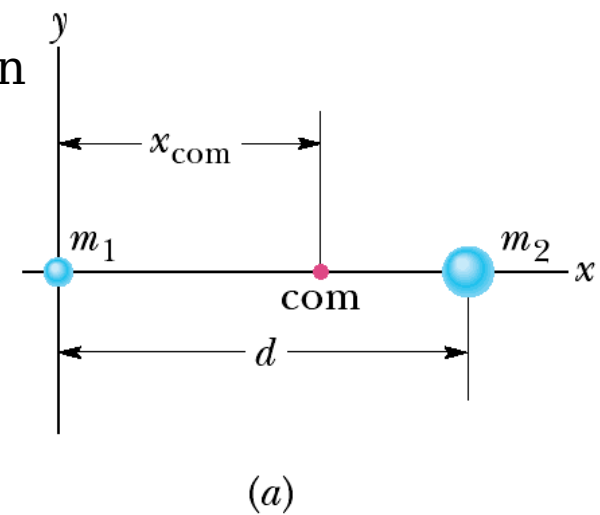
- extend this equation to a more general situation in which  $n$  particles are strung out along the  $x$  axis, then the location of the center of mass is

$$x_{\text{com}} = \frac{m_1 x_1 + m_2 x_2 + \dots + m_n x_n}{M} = \frac{1}{M} \sum_{i=1}^n m_i x_i$$

where  $M = m_1 + m_2 + \dots + m_n$

- In 3-D situations, the center of mass is identified by 3 coordinates as

$$x_{\text{com}} = \frac{1}{M} \sum_{i=1}^n m_i x_i, \quad y_{\text{com}} = \frac{1}{M} \sum_{i=1}^n m_i y_i, \quad z_{\text{com}} = \frac{1}{M} \sum_{i=1}^n m_i z_i$$



- With the language of vector, the position of a particle is  $\vec{r}_i = x_i \hat{i} + y_i \hat{j} + z_i \hat{k}$  then the position of the center of mass of a system of particles is

$$\vec{r}_{\text{com}} = x_{\text{com}} \hat{i} + y_{\text{com}} \hat{j} + z_{\text{com}} \hat{k}$$

combining the equations above gives

$$\vec{r}_{\text{com}} = \frac{1}{M} \sum_{i=1}^n m_i \vec{r}_i$$

## Solid Bodies

- An ordinary object contains so many particles that we can best treat it as a continuous distribution of matter.

- the coordinates of the center of mass are defined as

$$x_{\text{com}} = \frac{1}{M} \int x \, d m, \quad y_{\text{com}} = \frac{1}{M} \int y \, d m, \quad z_{\text{com}} = \frac{1}{M} \int z \, d m$$

- For a uniform object, the density, or mass per unit volume, is uniform,

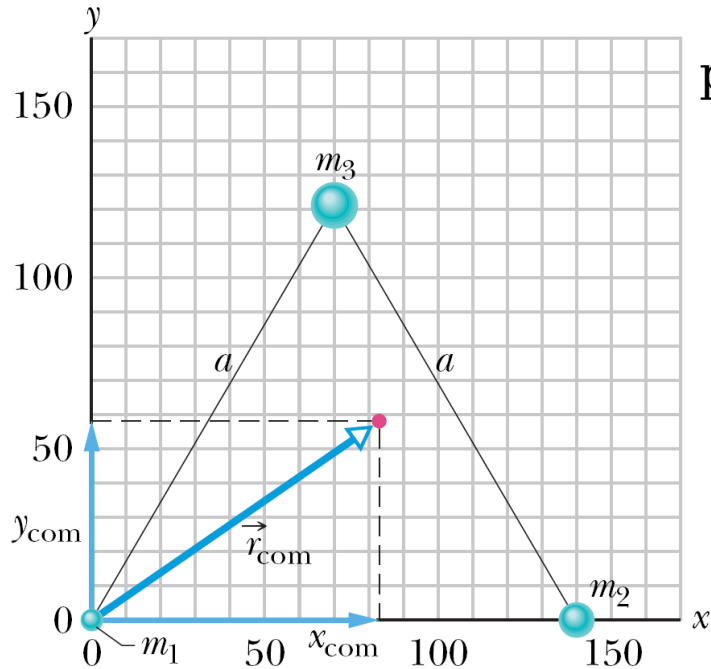
$$\rho = \frac{d m}{d V} = \frac{M}{V}$$

Using  $dm = (M/V) dV$ ,

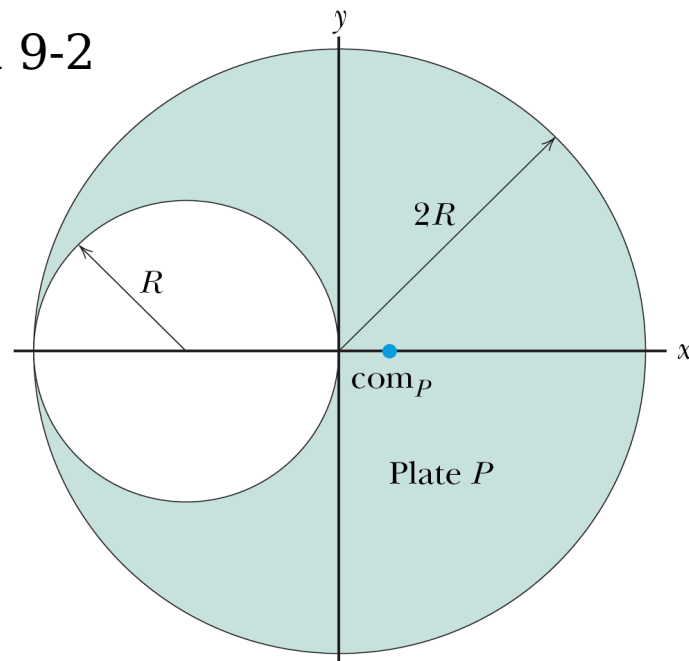
$$x_{\text{com}} = \frac{1}{V} \int x \, d V, \quad y_{\text{com}} = \frac{1}{V} \int y \, d V, \quad z_{\text{com}} = \frac{1}{V} \int z \, d V$$

- The center of mass of an object need not lie within the object.

problem 9-1



problem 9-2

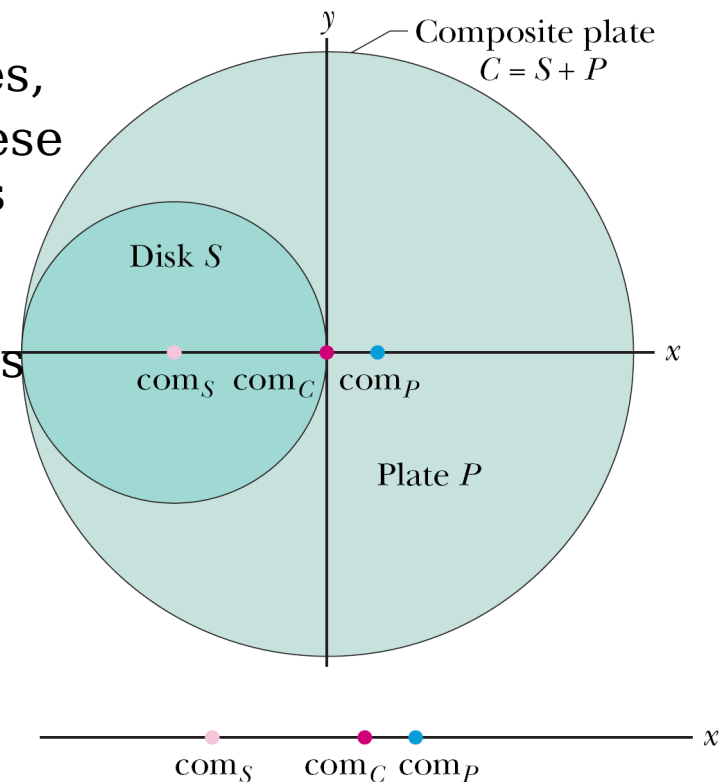


## Newton's 2<sup>nd</sup> Law for a System of Particles

- For an assemblage of  $n$  particles of different masses, we are interested not in the individual motions of these particles but only in the motion of the center of mass of the assemblage.

- Although the center of mass is just a point, it moves like a particle whose mass is equal to the total mass of the system; we can assign a position, a velocity, and an acceleration to it.

- The Newton's 2<sup>nd</sup> law for the motion of the center of mass of a system



$$\vec{F}_{\text{net}} = M \vec{a}_{\text{com}} \quad \text{system of particles}$$

- (1)  $\vec{F}_{\text{net}}$  is the net force of *all external forces* that act on the system;
- (2)  $M$  is the *total mass* of the system;
- (3)  $\vec{a}_{\text{com}}$  is the acceleration of the *center of mass* of the system.

● We can decompose the above equation into

$$F_{\text{net}, x} = M a_{\text{com}, x}, \quad F_{\text{net}, y} = M a_{\text{com}, y}, \quad F_{\text{net}, z} = M a_{\text{com}, z}$$

### Proof of the Equation

For a system of  $n$  particles:  $M \vec{r}_{\text{com}} = m_1 \vec{r}_1 + m_2 \vec{r}_2 + m_3 \vec{r}_3 + \dots + m_n \vec{r}_n$

Differentiating the above equation with respect time gives

$$M \vec{v}_{\text{com}} = m_1 \vec{v}_1 + m_2 \vec{v}_2 + m_3 \vec{v}_3 + \dots + m_n \vec{v}_n$$

Differentiating the above equation with respect time gives

$$M \vec{a}_{\text{com}} = m_1 \vec{a}_1 + m_2 \vec{a}_2 + m_3 \vec{a}_3 + \dots + m_n \vec{a}_n$$

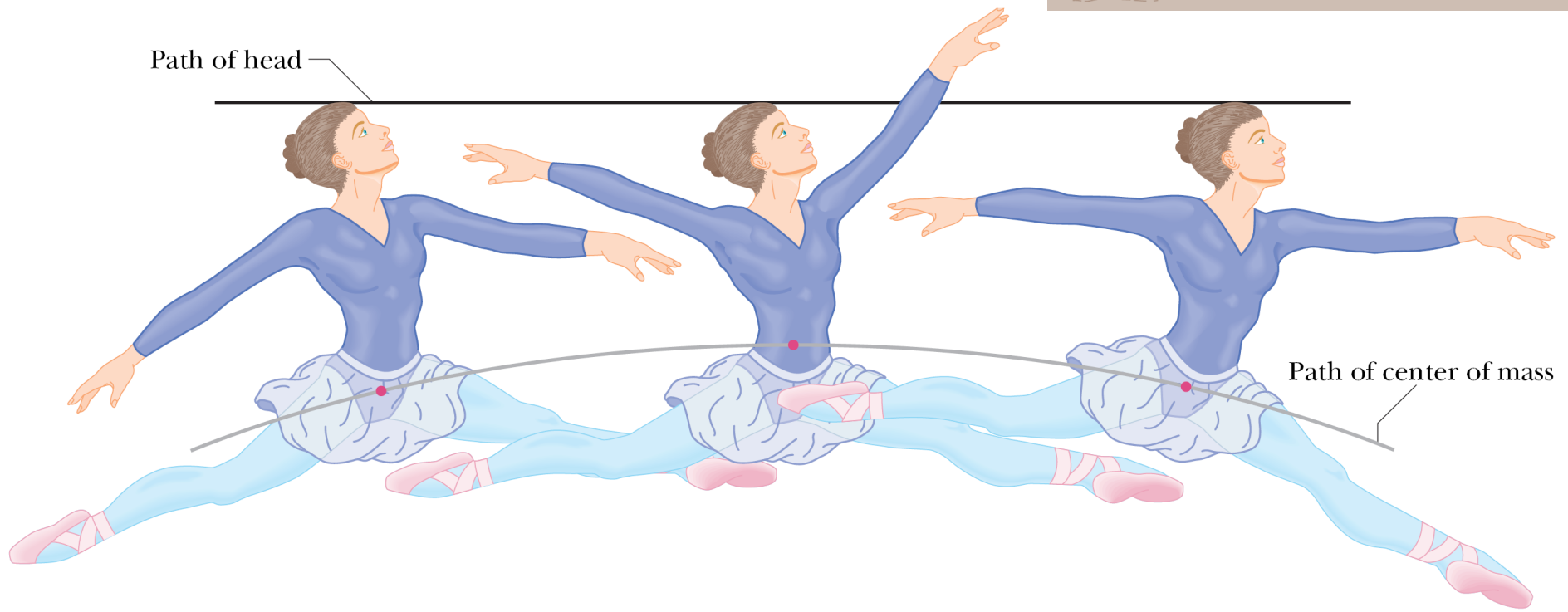
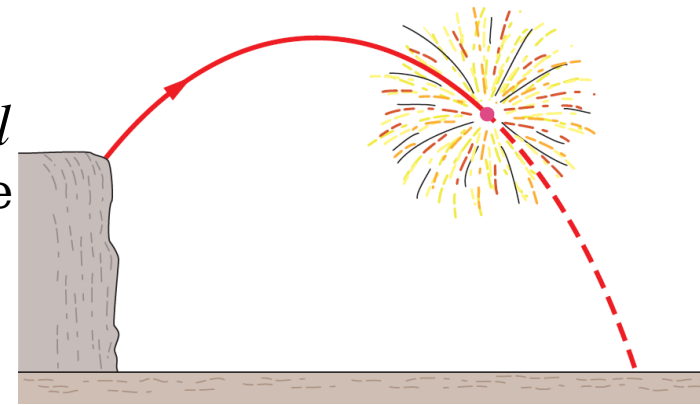
Although the center of mass is just a geometrical point, it has a position, a velocity, and an acceleration, as if it were a particle.

From Newton's 2<sup>nd</sup> law,  $M \vec{a}_{\text{com}} = \vec{F}_1 + \vec{F}_2 + \vec{F}_3 + \dots + \vec{F}_n$

- the forces in the right side will be either the internal forces or the external forces.
- By Newton's 3<sup>rd</sup> law, the internal forces form third-law force pairs and cancel out in the sum.
- The left is the vector sum of all the external forces that act on the system, thus

$$\vec{F}_{\text{net}} = M \vec{a}_{\text{com}}$$

- The forces of the explosion of a rocket are *internal* to the system; that is, they are forces on parts of the system from other parts.



- The net *external* force acting on the system is the gravitational force on the system, regardless of whether the rocket explodes. Thus,  $\vec{a}_{\text{com}} = \vec{g}$
- The center of mass of the fragments follows the same parabolic trajectory that the rocket would have followed had it not exploded.
- Although a ballet dancer raises her arms and stretches her legs out to shift her center of mass upward through her body, the center of mass still follows a parabolic path across the stage.

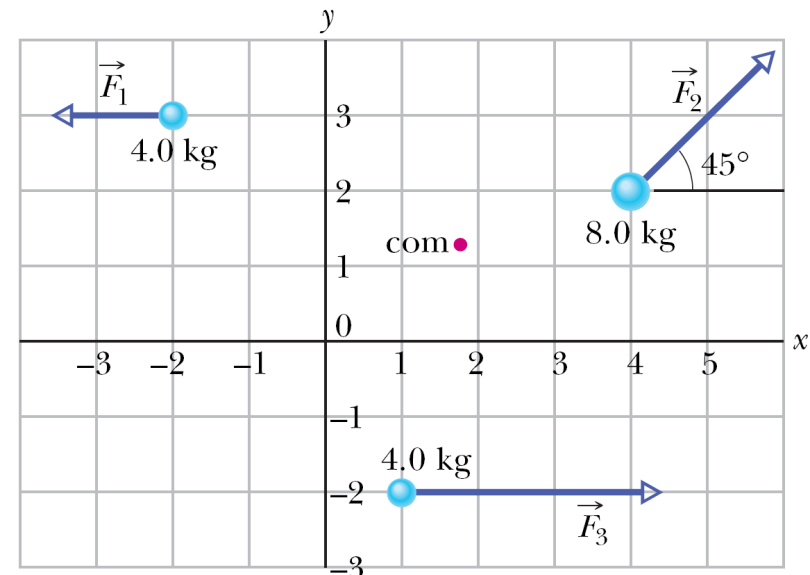
problem 9-3

## Linear Momentum

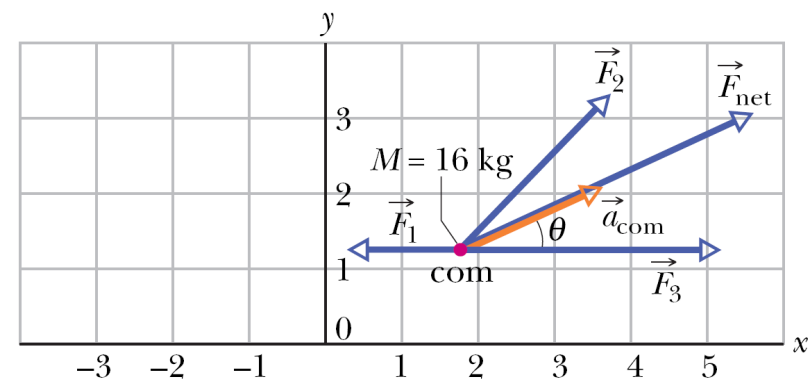
- The **linear momentum** of a particle is a vector quantity that is defined as

$$\vec{p} = m \vec{v} \quad \text{linear momentum of a particle}$$

- Since  $m$  is always a positive scalar quantity, the momentum and the velocity have the same direction.
- the SI unit for momentum is the kilogram-meter per second ( $\text{kg} \cdot \text{m/s}$ ).



(a)



(b)

- expressed Newton's 2<sup>nd</sup> law of motion in terms of momentum:

The time rate of change of the momentum of a particle is equal to the net force acting on the particle and is in the direction of that force.

In equation form this becomes

$$\vec{F}_{\text{net}} = \frac{d \vec{p}}{d t}$$

- the linear momentum can be changed only by a net external force. If there is no net external force, the momentum *cannot* change.

- the equation can be manipulated as  $\vec{F}_{\text{net}} = \frac{d \vec{p}}{d t} = \frac{d}{d t} (m \vec{v}) = m \frac{d \vec{v}}{d t} = m \vec{a}$

- the relations  $\vec{F}_{\text{net}} = d \vec{p} / d t$  and  $\vec{F}_{\text{net}} = m \vec{a}$  are equivalent expressions of Newton's 2<sup>nd</sup> law of motion for a particle.

## The Linear Momentum of a System of Particles

- Consider a system of  $n$  particles, each with its own mass, velocity, and linear momentum. The system as a whole has a total linear momentum, which is defined to be the vector sum of the individual particles linear momenta,

$$\begin{aligned} \vec{P} &= \vec{p}_1 + \vec{p}_2 + \vec{p}_3 + \cdots + \vec{p}_n \\ &= m_1 \vec{v}_1 + m_2 \vec{v}_2 + m_3 \vec{v}_3 + \cdots + m_n \vec{v}_n \end{aligned}$$

thus,

$$\vec{P} = M \vec{v}_{\text{com}} \quad \text{linear momentum, system of particles}$$

The linear momentum of a system of particles is equal to the product of the total mass  $M$  of the system and the velocity of the center of mass.

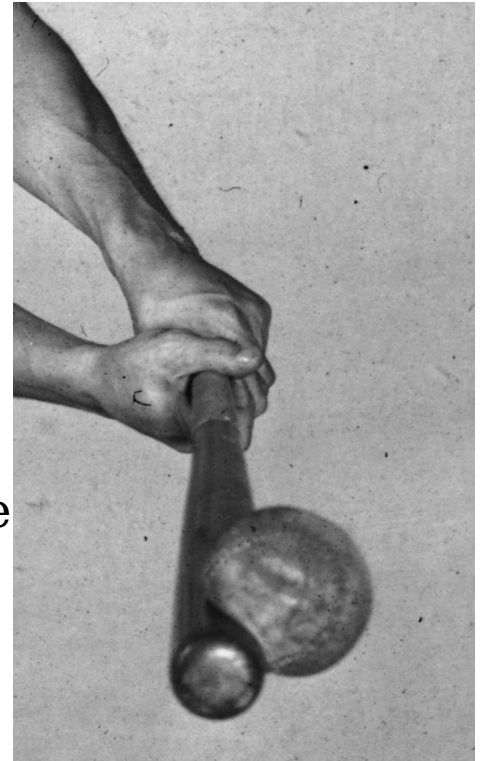
- Taking the time derivative of the above equation gives

$$\frac{d \vec{P}}{d t} = M \frac{d \vec{v}_{\text{com}}}{d t} = M \vec{a}_{\text{com}}$$

- thus, Newton's 2<sup>nd</sup> law for a system of particles is

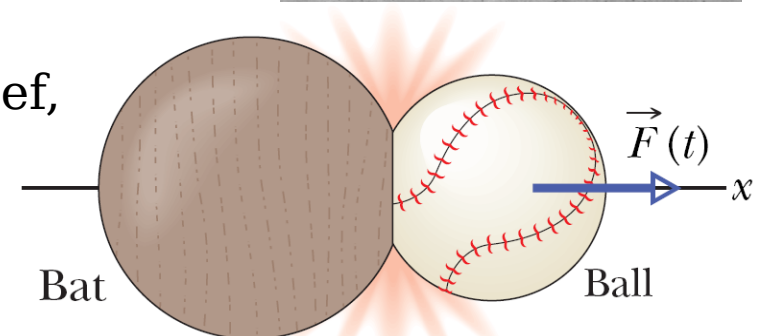
$$\vec{F}_{\text{net}} = \frac{d \vec{P}}{d t} \quad \text{system of particles}$$

- The net external force on a system of particles changes the linear momentum of the system. And the linear momentum can be changed only by a net external force. If there is no net external force, the linear momentum *cannot* change.



## Collision and Impulse

- In a collision, the external force on a body is brief, has large magnitude, and suddenly changes the body's momentum.



## Single Collision

- A object experiences a force that varies during the collision and changes the linear momentum of the object. Thus, in time interval  $dt$ , the change in the object's momentum is

$$d \vec{p} = \vec{F}(t) d t$$

- The net change in the object's momentum due to the collision from a time  $t_i$  to  $t_f$  is

$$\vec{J} \leftarrow \Delta \vec{p} \leftarrow \vec{p}_f - \vec{p}_i \leftarrow \int_{t_i}^{t_f} d\vec{p} = \int_{t_i}^{t_f} \vec{F}(t) dt$$

- The impulse of the collision is a measure of both the magnitude and the duration of the collision force:

$$\vec{J} = \int_{t_i}^{t_f} \vec{F}(t) dt \quad \text{impulse defined}$$

- the change in an object's momentum is equal to the impulse on the object:

$$\Delta \vec{p} = \vec{J} \quad \text{linear momentum-impulse theorem}$$

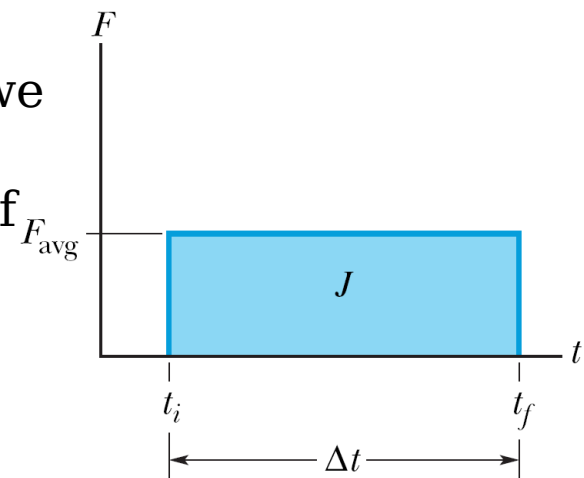
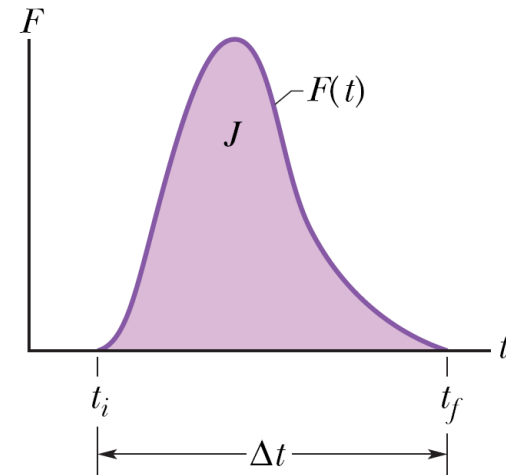
- In a component form, eg, for  $x$  axis,

$$J_x \leftarrow \Delta p_x \leftarrow p_{fx} - p_{ix} \leftarrow \int_{t_i}^{t_f} dp_x = \int_{t_i}^{t_f} F_x dt$$

- If the function of the force is known, we can evaluate the impulse by integrating the force function.

- In many situations, the force function is unknown, but we do know the average magnitude  $F_{\text{avg}}$  of the force and the duration  $\Delta t$  ( $= t_f - t_i$ ) of the collision, then the magnitude of the impulse can be

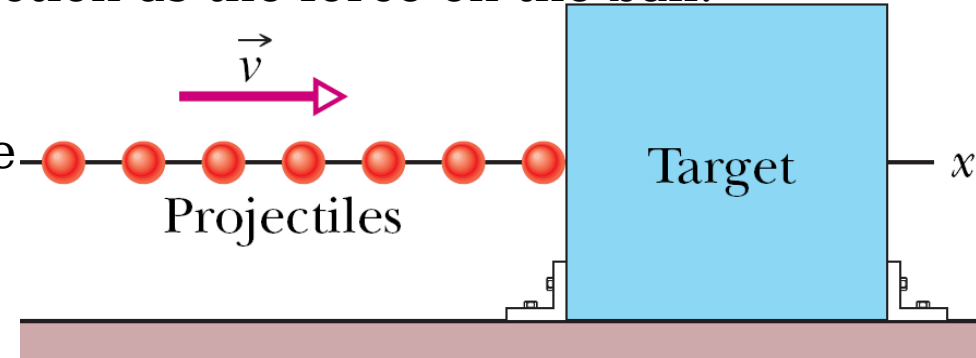
$$J = F_{\text{avg}} \Delta t$$



- For a bat hitting a ball, Newton's 3<sup>rd</sup> law tells us that the force on the bat has the same magnitude but the opposite direction as the force on the ball.

### Series of Collisions

- We would like to know the average force during a large number of collisions.



- Let  $n$  be the number of projectiles that collide in a time interval  $\Delta t$ . Thus each projectile has initial momentum  $mv$  and undergoes a change  $\Delta p$  in linear momentum because of the collision.

- The total change in linear momentum for  $n$  projectiles during interval  $\Delta t$  is  $n\Delta p$ , thus the impulse on the target is  $J = -n \Delta p$

the minus sign indicates that  $J$  and  $\Delta p$  have opposite directions.

- The average force acting on the target is  $F_{\text{avg}} = \frac{J}{\Delta t} = \frac{-n}{\Delta t} \Delta p = \frac{-n}{\Delta t} m \Delta v$

- If the projectiles stop upon impact, then  $\Delta v = v_f - v_i = 0 - v = -v$

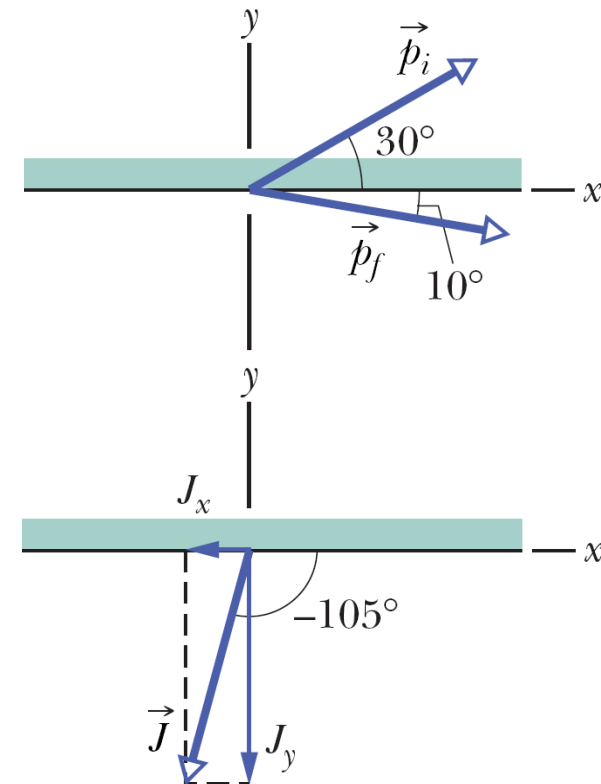
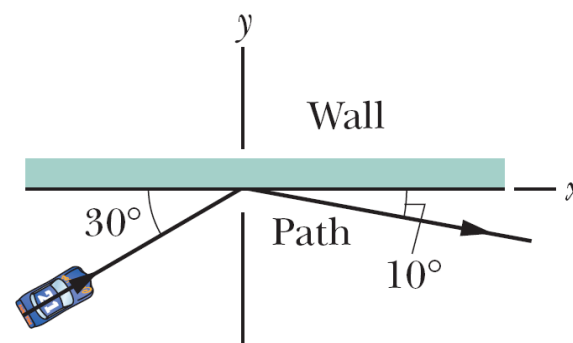
- If the projectiles bounce directly backward from the target with no change in speed, then

$$\Delta v = v_f - v_i = -v - v = -2v$$

- In time interval  $\Delta t$ , an amount of mass  $\Delta m = nm$  collides with the target, then we can rewrite the equation

$$F_{\text{avg}} = - \frac{\Delta m}{\Delta t} \Delta v$$

problem 9-4



## Conservation of Linear Momentum

- If the net external force (and thus the net impulse) acting on a system of particles is 0 (the system is isolated) and no particles leave or enter the system (the system is closed), then

$$\frac{d\vec{P}}{dt} = 0 \Rightarrow \vec{P} = \text{constant} \quad \text{closed, isolated system}$$

If no net external force acts on a system of particles, the total linear momentum of the system cannot change.

This result is called the **law of conservation of linear momentum**.

- The law can also be written as  $\vec{p}_i = \vec{p}_f$  closed, isolated system and

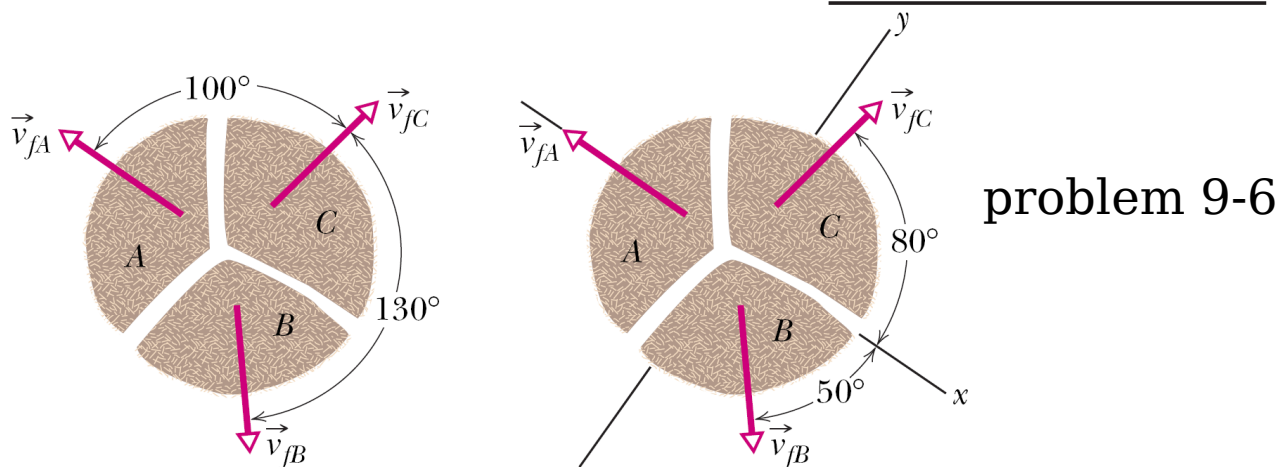
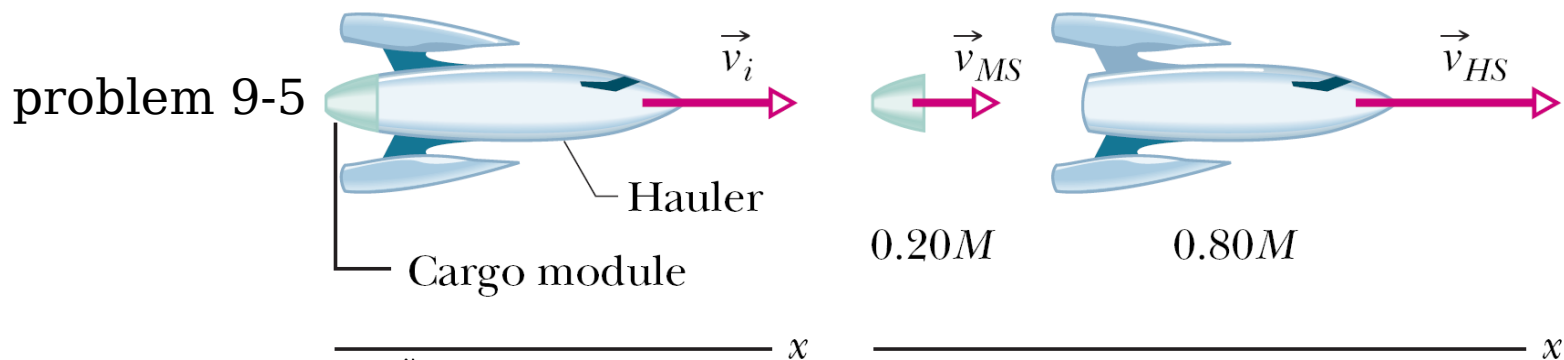
$$\left( \begin{array}{c} \text{total linear momentum} \\ \text{at some initial time } t_i \end{array} \right) = \left( \begin{array}{c} \text{total linear momentum} \\ \text{at some later time } t_f \end{array} \right)$$

- *Caution:* Momentum should not be confused with energy. In some cases, momentum is conserved but energy is definitely not.

- In components

If the component of the net *external* force on a closed system is 0 along an axis, then the component of the linear momentum of the system along that axis cannot change.

- Although internal forces can change the linear momentum of portions of the system, they cannot change the total linear momentum of the entire system.



# Momentum and Kinetic Energy in Collisions

● Assume the system is closed and isolated, then the law of conservation of linear momentum is a very powerful rule because it can allow us to determine the results of a collision *without* knowing the details of the collision.

● For the total kinetic energy of a system of 2 colliding bodies,

**elastic collision:** the kinetic energy of the system is *conserved* during the collision.

**inelastic collision:** the kinetic energy of the system is *not* conserved during the collision.

● In an inelastic collision, some energy is transferred from kinetic energy to other forms of energy, eg, thermal energy, energy of sound, etc.

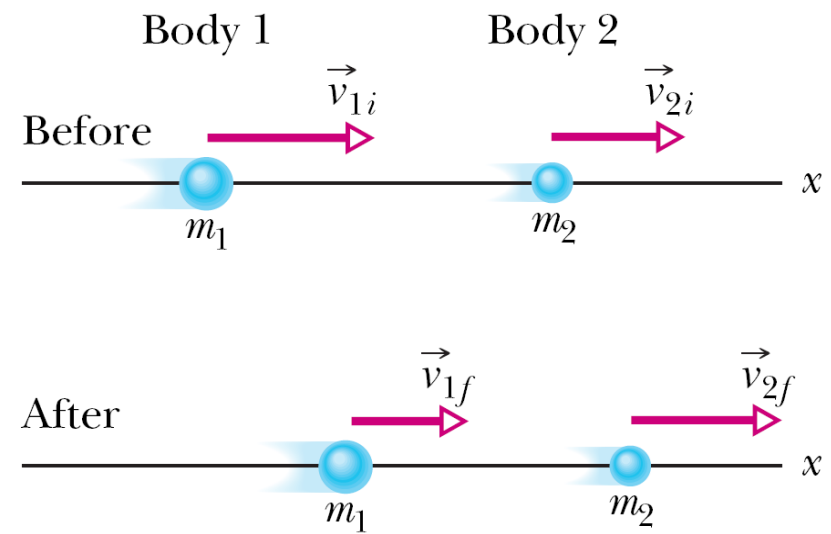
● In an inelastic collision, the greatest loss occurs if the bodies stick together, in which case the collision is called a **completely inelastic collision**.

## Inelastic Collisions in One Dimension

### One-Dimensional Inelastic Collision

● the law of conservation of linear momentum for this 2-body system is

$$\left( \begin{array}{l} \text{total momentum } \vec{P}_i \\ \text{before the collision} \end{array} \right) = \left( \begin{array}{l} \text{total momentum } \vec{P}_f \\ \text{after the collision} \end{array} \right)$$



$$\vec{p}_{1i} + \vec{p}_{2i} = \vec{p}_{1f} + \vec{p}_{2f} \quad \text{conservation of linear momentum}$$

Since the motion is 1-dim, we can rewrite the equation as

$$m_1 v_{1i} + m_2 v_{2i} = m_1 v_{1f} + m_2 v_{2f}$$

- If we know values for, say, the masses, the initial velocities, and one of the final velocities, we can find the other final velocity with this equation.

### One-Dimensional Completely Inelastic Collision

- For the body with mass  $m_2$  being at rest ( $v_{2i}=0$ ), we refer to that body as the *target* and to the incoming body as the *projectile*, and

$$m_1 v_{1i} = (m_1 + m_2) V$$

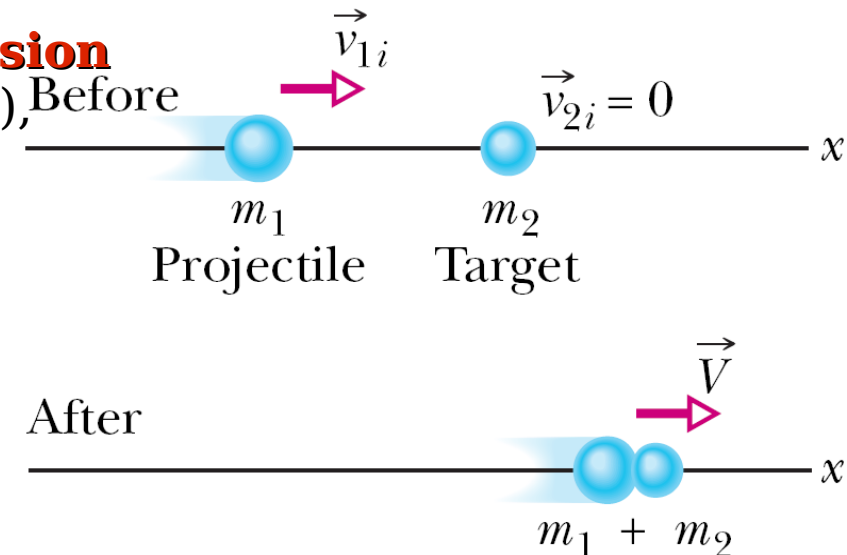
or

$$V = \frac{m_1}{m_1 + m_2} v_{1i}$$

- Note that  $V$  must be less than  $v_{1i}$  because the mass ratio  $m_1/(m_1+m_2)$  must be less than unity.

### Velocity of the Center of Mass

- In a closed, isolated system, the velocity of the center of mass of the system cannot be changed by a collision because there is no net external force to change it.



- relate the velocity to the total linear momentum by

$$\vec{P} = M \vec{v}_{\text{com}} = (m_1 + m_2) \vec{v}_{\text{com}}$$

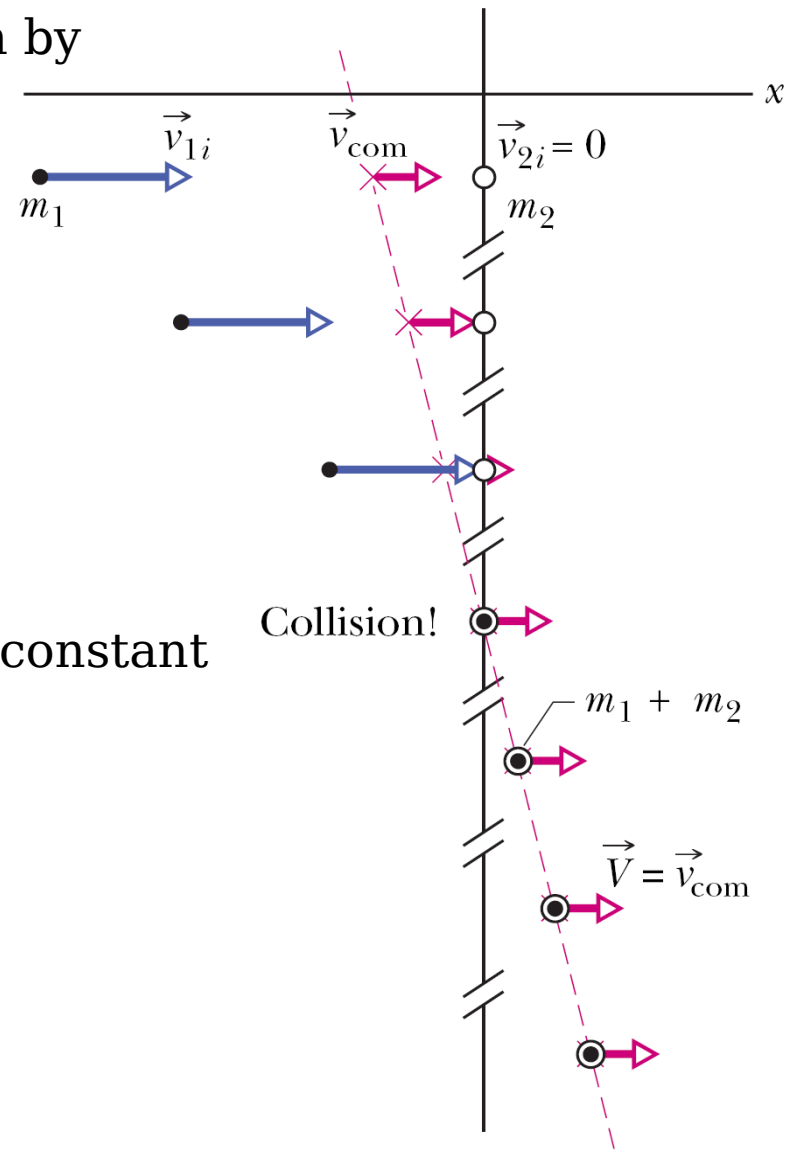
- The total linear momentum is conserved during the collision; and it is given from LHS

$$\vec{P} = \vec{p}_{1i} + \vec{p}_{2i}$$

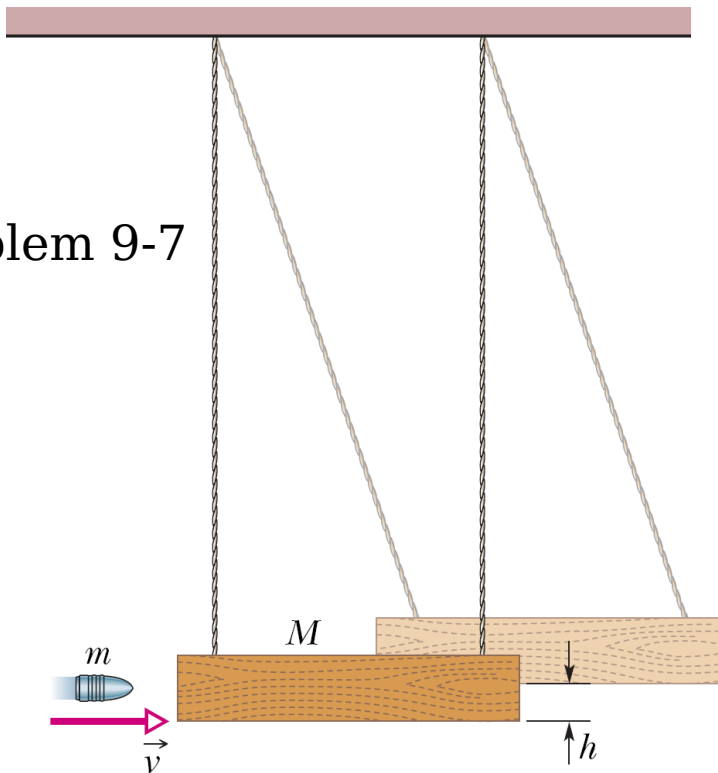
then

$$\vec{v}_{\text{com}} = \frac{\vec{P}}{m_1 + m_2} = \frac{\vec{p}_{1i} + \vec{p}_{2i}}{m_1 + m_2}$$

- the velocity of the center of mass has the same constant value before and after the collision.



problem 9-7



## Elastic Collisions in One Dimension

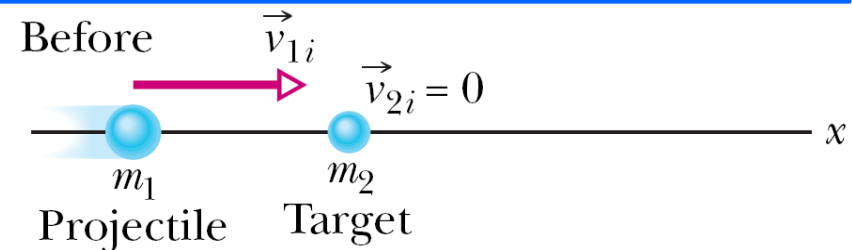
- In an elastic collision,  $\left( \begin{array}{l} \text{total kinetic energy} \\ \text{before the collision} \end{array} \right) = \left( \begin{array}{l} \text{total kinetic energy} \\ \text{after the collision} \end{array} \right)$

and it means

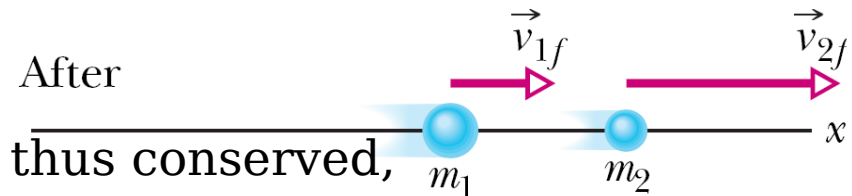
In an elastic collision, the kinetic energy of each colliding body may change, but the total kinetic energy of the system does not change.

### Stationary Target

- A projectile body of mass  $m_1$  and initial velocity  $v_{1i}$  moves toward a target body of mass  $m_2$  that is initially at rest ( $v_{2i} = 0$ ).



- If the 2-body system is closed and isolated, and the net linear momentum of the system is thus conserved,



$$m_1 v_{1i} = m_1 v_{1f} + m_2 v_{2f} \quad \text{linear momentum}$$

- If the collision is also elastic, then the total kinetic energy is conserved,

$$\frac{1}{2} m_1 v_{1i}^2 = \frac{1}{2} m_1 v_{1f}^2 + \frac{1}{2} m_2 v_{2f}^2 \quad \text{kinetic energy}$$

- If we know the masses of the bodies and if we also know  $v_{1i}$ , the only unknown quantities are  $v_{1f}$  and  $v_{2f}$ , the final velocities of the 2 bodies, then

$$m_1 (v_{1i} - v_{1f}) = m_2 v_{2f}$$

and

$$m_1 (v_{1i} - v_{1f})(v_{1i} + v_{1f}) = m_2 v_{2f}^2$$

we can obtain

$$v_{1f} = \frac{m_1 - m_2}{m_1 + m_2} v_{1i}$$
$$v_{2f} = \frac{2 m_1}{m_1 + m_2} v_{1i}$$

●  $v_{2f}$  is always positive (the initially stationary target body with mass  $m_2$  always moves forward).

●  $v_{1f}$  is positive if  $m_1 > m_2$ , is negative if  $m_1 < m_2$ .

● some special cases:

1. **Equal masses**: ie,  $m_1 = m_2$ , then  $v_{1f} = 0$  and  $v_{2f} = v_{1i}$ .

this means in head-on collisions, bodies of equal mass simply exchange velocities. This is true even if body 2 is not initially at rest.

2. **A massive target**: ie,  $m_1 \ll m_2$ , then  $v_{1f} \approx -v_{1i}$  and  $v_{2f} \approx \frac{2 m_1}{m_2} v_{1i}$

- It tells that the projectile simply bounces back along its incoming path, its speed essentially unchanged. And the target moves forward at a low speed.

3. **A massive projectile**: ie,  $m_1 \gg m_2$ , then  $v_{1f} \approx v_{1i}$  and  $v_{2f} \approx 2 v_{1i}$

- It tells that the projectile simply keep on going, scarcely slowed by the collision; the target moves forward at twice speed of the projectile.

### Moving target

- examine the situation in which both bodies are moving before they undergo an elastic collision.



- the conservation of linear momentum gives  $m_1 v_{1i} + m_2 v_{2i} = m_1 v_{1f} + m_2 v_{2f}$

- the conservation of kinetic energy gives  $\frac{1}{2} m_1 v_{1i}^2 + \frac{1}{2} m_2 v_{2i}^2 = \frac{1}{2} m_1 v_{1f}^2 + \frac{1}{2} m_2 v_{2f}^2$

- To solve these equation for  $v_{1f}$  and  $v_{2f}$

$$m_1 (v_{1i} - v_{1f}) = -m_2 (v_{2i} - v_{2f})$$

$$m_1 (v_{1i} - v_{1f}) (v_{1i} + v_{1f}) = -m_2 (v_{2i} - v_{2f}) (v_{2i} + v_{2f})$$

then we obtain

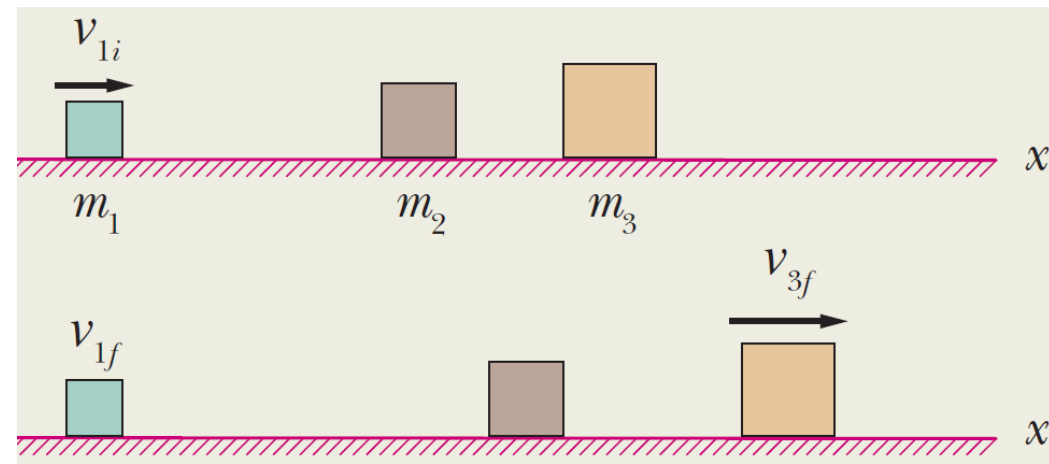
$$v_{1f} = \frac{m_1 - m_2}{m_1 + m_2} v_{1i} + \frac{2 m_2}{m_1 + m_2} v_{2i}$$

$$v_{2f} = \frac{2 m_1}{m_1 + m_2} v_{1i} + \frac{m_2 - m_1}{m_1 + m_2} v_{2i}$$

problem 9-8

## Collisions in Two Dimensions

- When a collision is not head-on, the bodies do not end up traveling along their initial axis.



- For a 2-dim collisions in a closed, isolated system, the conservation of linear momentum gives

$$\vec{P}_{1i} + \vec{P}_{2i} = \vec{P}_{1f} + \vec{P}_{2f}$$

- If the collision is elastic, the conservation of kinetic energy gives

$$K_{1i} + K_{2i} = K_{1f} + K_{2f}$$

- For a *glancing collision* (not head-on) between a projectile and a target initially at rest, the component of the linear momentum along x-axis is

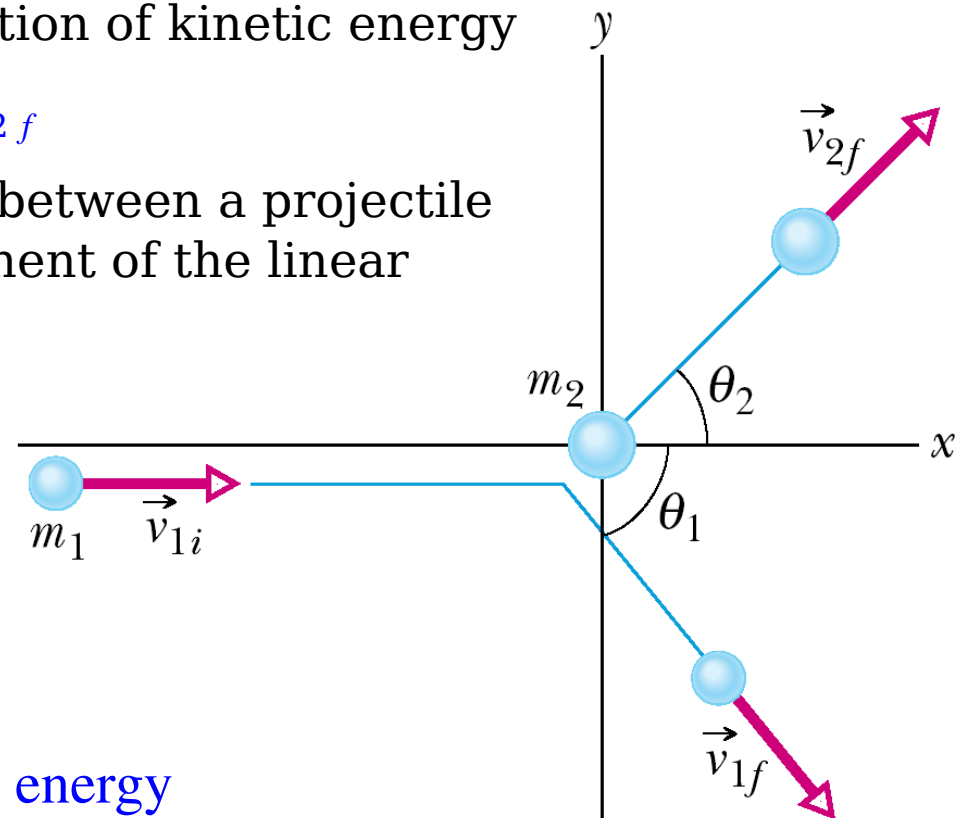
$$m_1 v_{1i} = m_1 v_{1f} \cos \theta_1 + m_2 v_{2f} \cos \theta_2$$

and along the y axis is

$$0 = -m_1 v_{1f} \sin \theta_1 + m_2 v_{2f} \sin \theta_2$$

- For an elastic collision,

$$\frac{1}{2} m_1 v_{1i}^2 = \frac{1}{2} m_1 v_{1f}^2 + \frac{1}{2} m_2 v_{2f}^2 \quad \text{kinetic energy}$$



- There are 7 variables:  $m_1, m_2, v_{1i}, v_{1f}, v_{2f}, \theta_1, \theta_2$ . If we know any 4 of these quantities, we can solve the 3 equations for the remaining 3 quantities.

## System with varying Mass: A Rocket

- We handle the variation of the mass of the rocket as the rocket accelerates by applying Newton's 2<sup>nd</sup> law, not to the rocket alone but to the rocket and its ejected combustion products taken together.

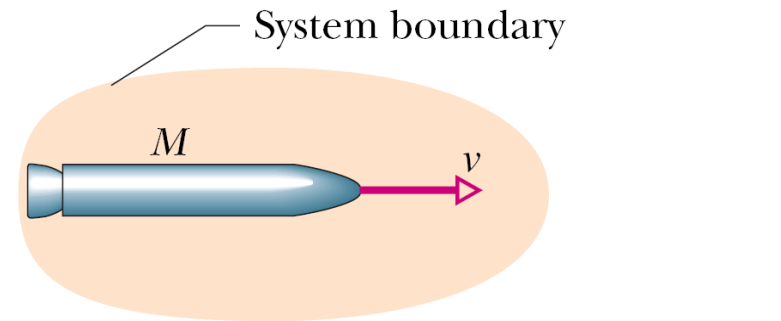
### Finding the Acceleration

- For the one-dimensional motion of a rocket, let  $M$  be the mass of the rocket and  $v$  its velocity at an arbitrary time  $t$ .

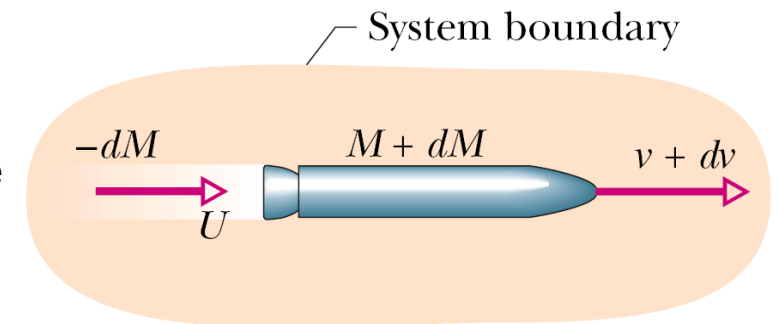
- After a time interval  $dt$ , the rocket has velocity  $v+dv$  and mass  $M+dM$ , where the change in mass  $dM$  is a negative quantity. The exhaust products has mass  $-dM$  and velocity  $U$ .

- The system, consisting of the rocket and its exhaust products, is closed and isolated, thus the linear momentum must be conserved during  $dt$ ,

$$P_i = P_f$$



(a)



(b)

Then

$$M v = -U dM + (M + dM)(v + dv)$$

where the 1<sup>st</sup> term in RHS is the linear momentum of the exhaust products and the 2<sup>nd</sup> term is the linear momentum of the rocket.

- We can simplify the equation by using the relative speed between the rocket and the exhaust products,

$$\left( \begin{array}{c} \text{velocity of rocket} \\ \text{relative to frame} \end{array} \right) = \left( \begin{array}{c} \text{velocity of rocket} \\ \text{relative to products} \end{array} \right) + \left( \begin{array}{c} \text{velocity of products} \\ \text{relative to frame} \end{array} \right)$$

That is,  $(v + dv) = v_{\text{rel}} + U$  or  $U = v + dv - v_{\text{rel}}$

This gives  $-v_{\text{rel}} dM = M dv \Rightarrow -\frac{dM}{dt} v_{\text{rel}} = M \frac{dv}{dt}$

- define  $R$  to be the (positive) mass rate of fuel consumption, then  $dM/dt = -R$ ,

$$R v_{\text{rel}} = M a \quad \text{1st rocket equation}$$

- the LHS of the equation has the dimension of force ( $\text{kg/s m/s} = \text{kg m/s}^2 = \text{N}$ ) and depends only on design characteristics of the rocket engine.

- We call the term  $Rv_{\text{rel}}$  the thrust of the rocket engine and represents it with  $T$ , then  $T = Ma$ .

## Finding the Velocity

● since

$$d v = -v_{\text{rel}} \frac{d M}{M} \Rightarrow \int_{v_i}^{v_f} d v = -v_{\text{rel}} \int_{M_i}^{M_f} \frac{d M}{M}$$

Evaluating the integrals gives

$$v_f - v_i = v_{\text{rel}} \ln \frac{M_i}{M_f} \quad \text{2nd rocket equation}$$

● The advantage of multistage rockets is that  $M_f$  is reduced by discarding successive stages when their fuel is depleted.

Problem 9.9

The chosen problems: 13, 44, 68