

As the discussion in sec. 8.1, the current density and magnetic induction below the conducting surface are

$$\mathbf{J} = \frac{(1-i)}{\delta} \hat{\mathbf{n}} \times \mathbf{H}_{\parallel} e^{-\frac{\xi}{\delta}(1-i)} e^{-i\omega t} \quad \mathbf{B}_c = \mu_c \mathbf{H}_c = \mu_c \mathbf{H}_{\parallel} e^{-\frac{\xi}{\delta}} e^{i\frac{\xi}{\delta}} e^{-i\omega t}$$

(ref: textbook (8.13) and (8.9).)

The time-averaged force is

$$\langle \mathbf{F} \rangle = \left\langle \int \mathbf{J} \times \mathbf{B}^* d^3x \right\rangle = \left\langle \iint \mathbf{J} \times \mathbf{B}^* d\xi da \right\rangle = \int da \underbrace{\int d\xi \langle \mathbf{J} \times \mathbf{B}^* \rangle}_{=\mathbf{f}}$$

The time-averaged force per unit area \mathbf{f} can be rewritten as

$$\begin{aligned} \mathbf{f} &= \int d\xi \langle \mathbf{J} \times \mathbf{B}^* \rangle = \int d\xi \frac{1}{2} \text{Re} [\mathbf{J} \times \mathbf{B}^*] \\ &= \frac{1}{2} \text{Re} \int d\xi \left(\frac{(1-i)}{\delta} \hat{\mathbf{n}} \times \mathbf{H}_{\parallel} e^{-\frac{\xi}{\delta}(1-i)} e^{-i\omega t} \right) \times \left(\mu_c \mathbf{H}_{\parallel} e^{-\frac{\xi}{\delta}} e^{i\frac{\xi}{\delta}} e^{-i\omega t} \right)^* \\ &= \text{Re} \left[\frac{\mu_c(1-i)}{2\delta} (\hat{\mathbf{n}} \times \mathbf{H}_{\parallel}) \times \mathbf{H}_{\parallel}^* \int d\xi \left(e^{-\frac{\xi}{\delta}(1-i)} e^{-i\omega t} e^{-\frac{\xi}{\delta}} e^{-i\frac{\xi}{\delta}} e^{+i\omega t} \right) \right] \\ &= \text{Re} \left[\frac{\mu_c(1-i)}{2\delta} \left(-\hat{\mathbf{n}}(\mathbf{H}_{\parallel}^* \cdot \mathbf{H}_{\parallel}) + \mathbf{H}_{\parallel} \underbrace{(\hat{\mathbf{n}} \cdot \mathbf{H}_{\parallel}^*)}_{=0} \right) \int_0^{\infty} d\xi e^{-\frac{2\xi}{\delta}} \right] \\ &= \text{Re} \left[-\frac{\mu_c(1-i)|\mathbf{H}_{\parallel}|^2}{2\delta} \hat{\mathbf{n}} \cdot \frac{\delta}{2} \right] \\ &= -\hat{\mathbf{n}} \frac{\mu_c}{4} |\mathbf{H}_{\parallel}|^2 \end{aligned}$$